



PIC solvers for Intense Beams in Synchrotrons

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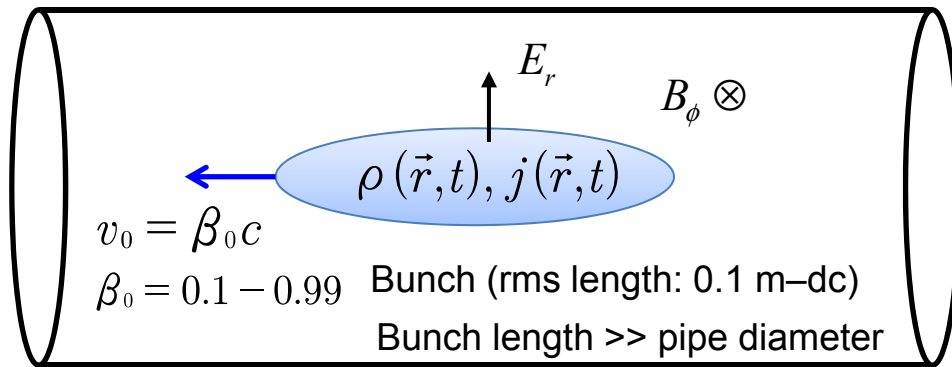
- Application and challenges for beam simulations with space charge
- 2D electrostatic Particle-In-Cell (PIC) simulation scheme for beams
- Review: Numerical noise in PIC codes and related emittance growth
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Space Charge Simulations in Synchrotrons

$$\epsilon_0 \nabla \cdot \vec{E} = \rho$$

(in the rest system of the beam)

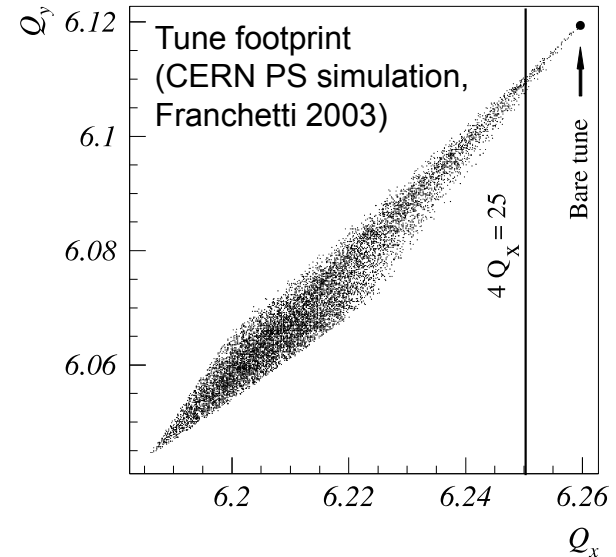
Space charge tune shift: $\Delta Q_y^{sc} \propto -\frac{q^2}{m} \frac{N}{B_f} \frac{4}{\epsilon \beta_0^2 \gamma_0^3}$



The transverse space charge force is the main intensity limiting effect in in the FAIR synchrotrons at GSI and in other high current synchrotrons.

Time scales: 1000-10⁶ turns (1 ms - 1 s)
 -> Emittance growth, Beam loss (< 1%)

Challenge: Control of numerical emittance growth !



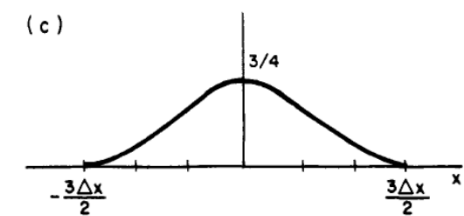
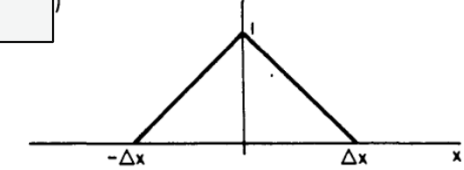
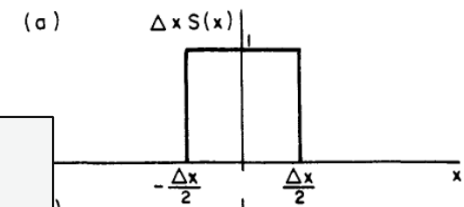
‘Space charge’ collaboration
 (CERN, GSI, FNAL, SNS, KEK,...)

- Codes for long-term simulations
 -> **Effect of stochastic noise !**
- (PIC) codes used for FAIR at GSI:
 pyORBIT, Micromap, PATRIC

PIC simulation scheme (for 2D beams)

2.5D electric fields:
 bunches are long (m)
 Compared to their
 transverse width (cm).

$$\rho(x, y, s) = Q \sum_{i=0}^M S(\vec{x} - \vec{x}_i)$$



q : beam particle charge

$Q = q \frac{N}{M}$: macro particle charge

N : number of beam particles

$M \ll N$: number of macro-particles

$$x_i'' - \kappa(s)x_i - \frac{qE_x(x_i, y_i, s)}{\gamma_0 m v_0^2} = 0$$

$$y_i'' + \kappa(s)y_i - \frac{qE_y(x_i, y_i, s)}{\gamma_0 m v_0^2} = 0$$

'Error' sources:

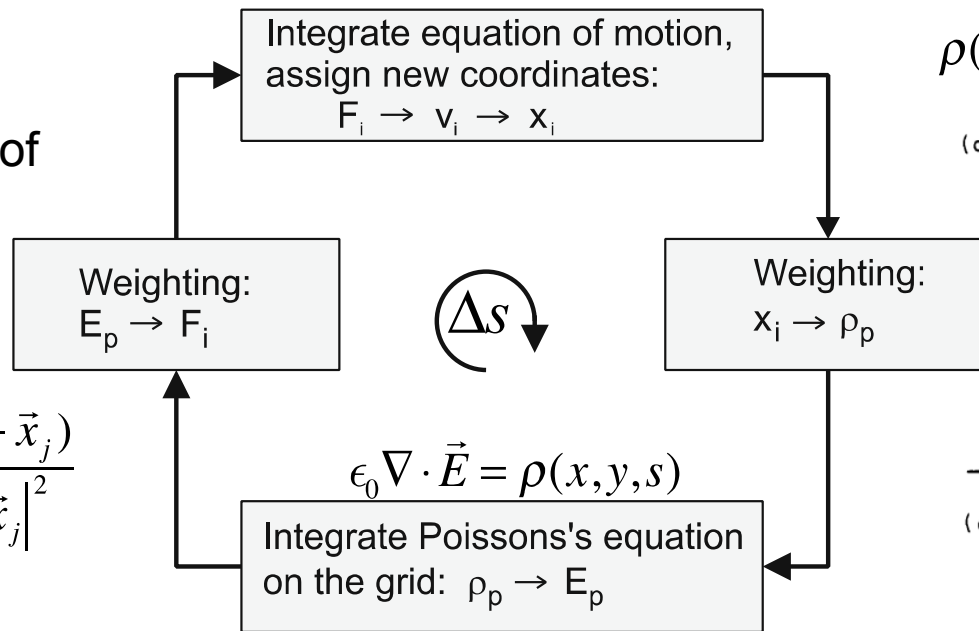
1) 'artificial' collisions of macro-particles Q and test particles q .

(Coulomb's law)

$$\vec{F}_i = \sum_{j=0}^M \frac{qQ'S(\vec{x}_i - \vec{x}_j)(\vec{x}_i - \vec{x}_j)}{2\pi\epsilon_0\gamma_0 m v_0^2 |\vec{x}_i - \vec{x}_j|^2}$$

-> energy conserving.

2) Grid/stochastic heating



$$\epsilon_0 \nabla \cdot \vec{E} = \rho(x, y, s)$$

Depending on particle shape S the macro-particles have a finite width: $\Delta \approx \Delta x$

What physics system does a PIC code represent ?

The system we would like to solve: **Collision-free Vlasov-Poisson (V-P) system**

$$\frac{\partial f}{\partial t} + \mathbf{x}' \frac{\partial f}{\partial \mathbf{x}} + \mathbf{F}_{\perp} \frac{\partial f}{\partial \mathbf{x}} = 0$$

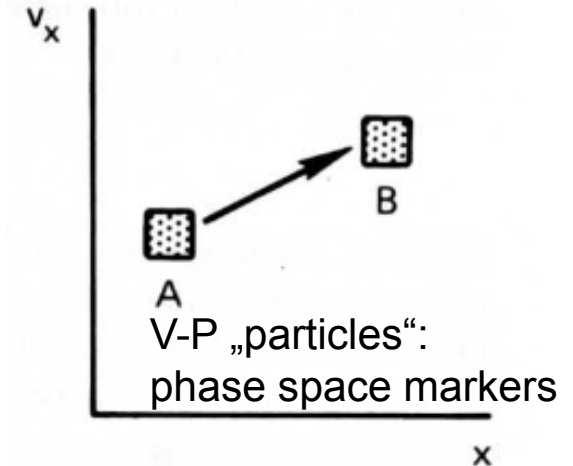
(Vlasov equation)

$$\epsilon_0 \nabla \phi = \rho(x, y)$$

(Poisson equation)

$$\rho(x, y) = q \int f(\mathbf{x}, \mathbf{x}') d^2 x \quad \mathbf{F}_{\perp} = \kappa_{x,y}(s) \mathbf{x} - \frac{q}{\gamma_0 m v_0^2} \nabla \phi(x, y, s)$$

(charge density) (AG focusing) (space charge)



$$\Rightarrow S = -k \int f \ln f dx dx' = \text{const.} \quad (\text{Entropy})$$

What in fact we solve is (forget about the grid for a moment):

$$\frac{\partial F}{\partial t} + \mathbf{x}' \frac{\partial F}{\partial \mathbf{x}} + \mathbf{F}_{\perp} \frac{\partial F}{\partial \mathbf{x}} = 0$$

(Klimontovich equation)

-> Artificial collisions and related noise.
-> Noise level is much stronger than the Schottky noise in real beams.

$$F(x, x') = \sum_{j=0}^M S(x - x_j) \delta(x' - x'_j) \quad \Rightarrow \quad S_{4M} = -k \int F \ln F dx dx' = \text{const.} \quad (\text{Entropy})$$

[1] Winjum et al., *Verification and convergence properties of PIC codes* (2014)

Artificial collisions in a 2D computer beam

In a 2D beam the beam macro-particles are rods: **Collision angle independent on b !**

All particles with relative velocities less than

$$v_{\perp}^2 = \frac{qQ'}{2\pi\epsilon_0 m} \text{ are deflected by angles } > 90^\circ.$$

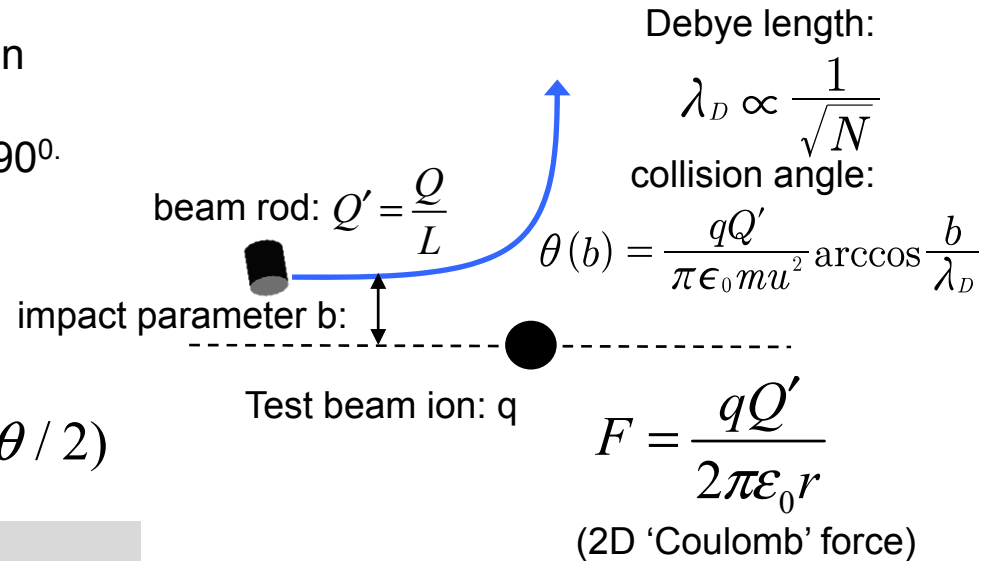
Friction force on a test particle:

$$F_p(\vec{v}) = m\vec{v}\vec{v} = m \int d^2v db u^2 f(\vec{v}) \sin^2(\theta/2)$$

$$v \propto \frac{N^2}{M} \left(1 - \frac{\Delta}{\lambda_D}\right) \quad \text{(2D collision rate for finite-sized macro particles)}$$

$$v \propto \frac{N^2}{M} \ln\left(\frac{\lambda_D}{\Delta}\right) \quad \text{(3D collision rate)}$$

$$\Delta \approx \Delta x \quad \text{(grid spacing)}$$



Relation between friction and diffusion in a system with energy conservation:

$$D = v \frac{k_B T}{m} \quad \text{(Einstein relation)}$$

$$T(s) = (T_x + T_y) / 2 \quad k_B T_x \propto mc^2 \frac{\epsilon_x}{\hat{\beta}_x}$$

Entropy and emittance growth due to artificial collisions in computer beams

J. Struckmeier, *Stochastic effects in real and simulated charged particle beams*, Phys. Rev. ST-AB 2000

Entropy/Emittance growth for anisotropic beams: $T_x \neq T_y$

$$D = v \frac{k_B T}{m} \Rightarrow \frac{1}{\varepsilon} \frac{d\varepsilon}{dt} = \frac{1}{2} k_B v \frac{(T_x - T_y)^2}{T_x T_y} \quad \varepsilon = \varepsilon_x \varepsilon_y \quad (4D \text{ emittance})$$

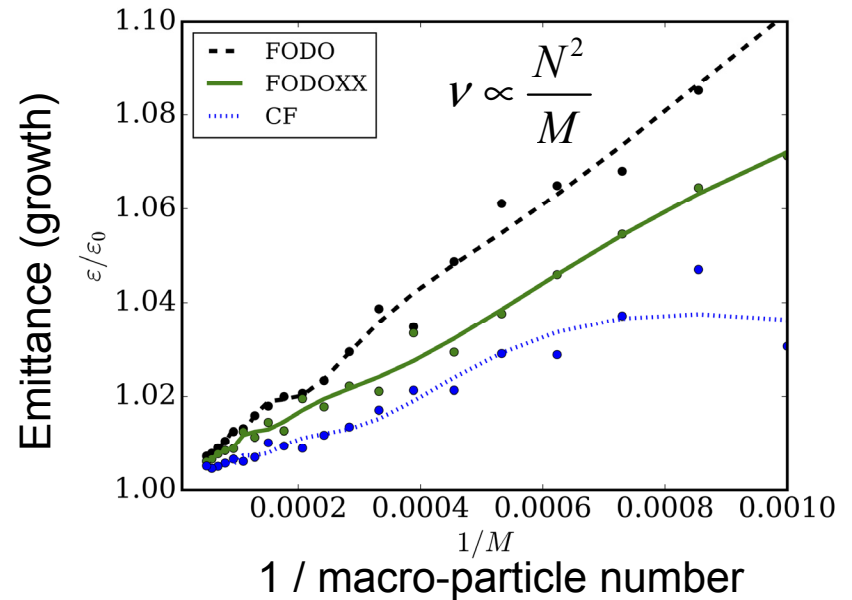
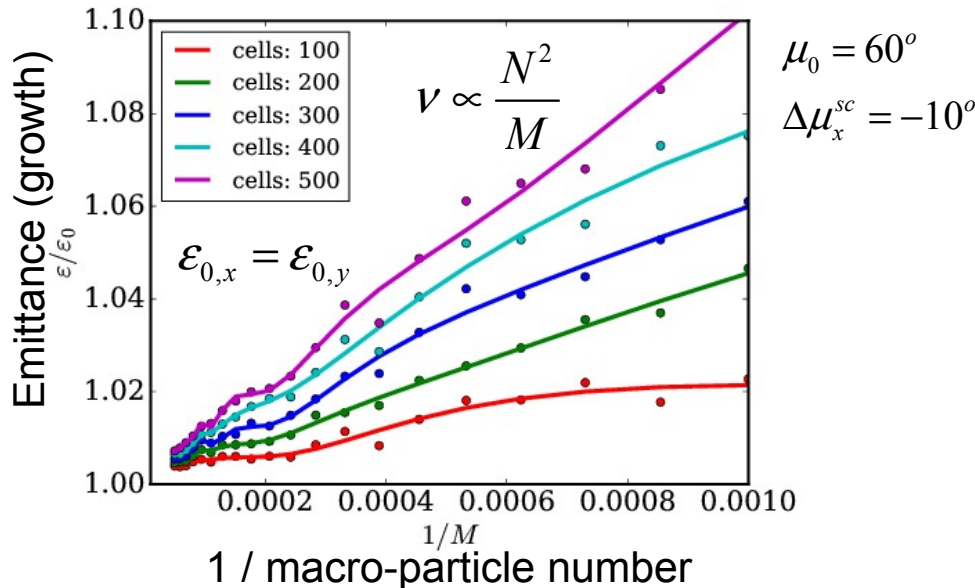
- > **Emittance growth due to AG focusing in a FODO channel !**
- > **No growth for constant focusing (CF) and for FODOXX channels.**

FODOXX: $\kappa_x(s) = \kappa_y(s) \Rightarrow \hat{\beta}_x = \hat{\beta}_y$

CF: $\kappa_x = \kappa_y = \text{const.}$

PATRIC (2D PIC):

Emittance growth after 100-500 FODO cells



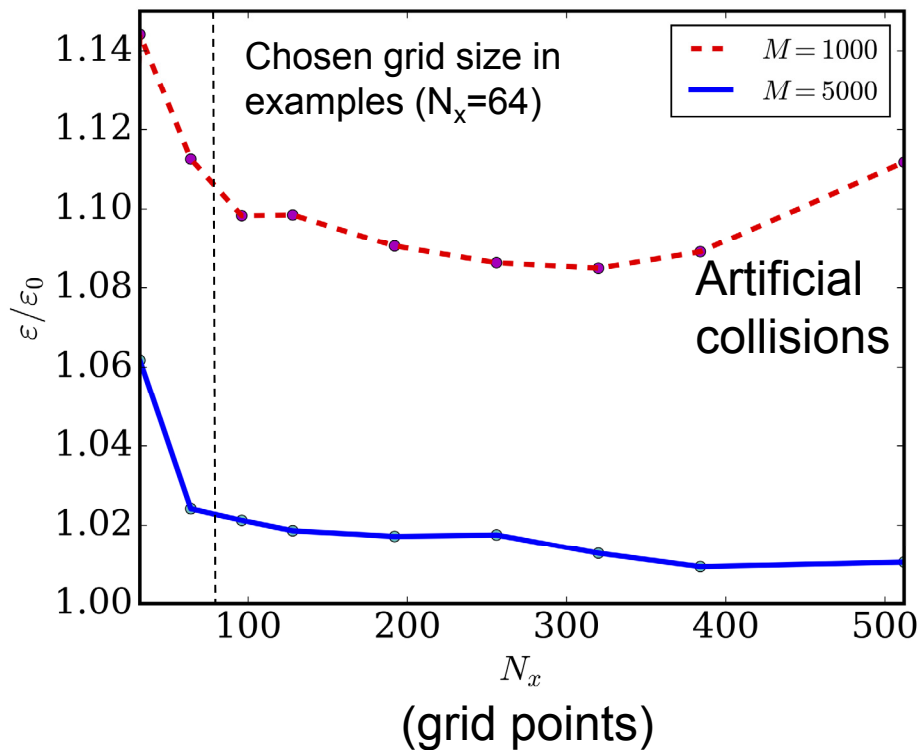
Emittance growth in CF/FODOXX channels.
-> **Obviously there is additional heating !**

[1] Hofmann, Boine-Frankenheim, PRAB (2015)

Short comment on the effect of the grid spacing on emittance growth

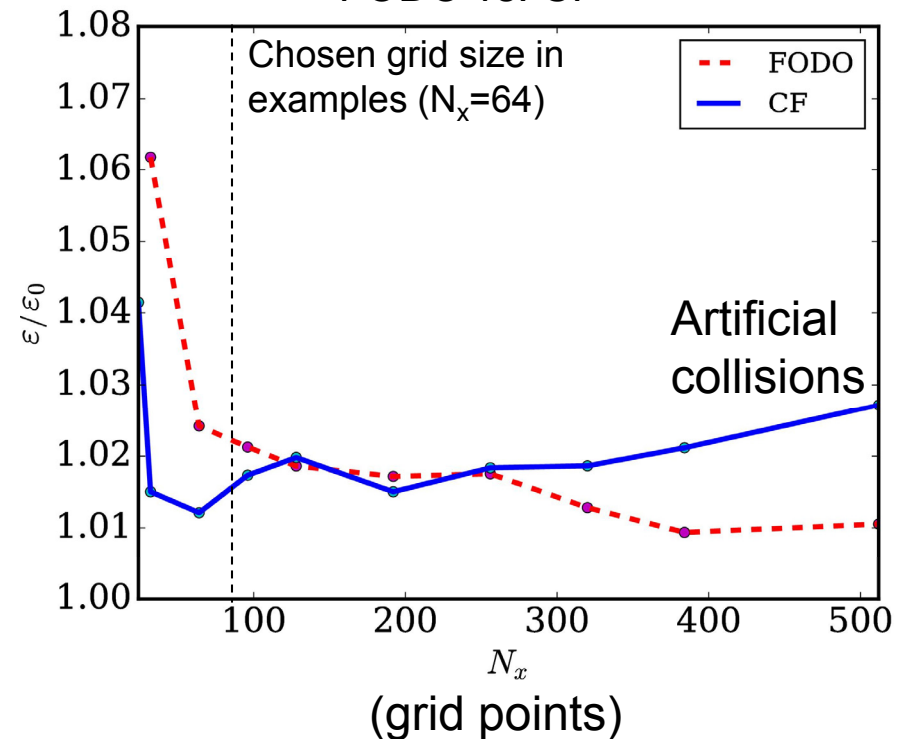
PATRIC (2D PIC):

Emittance growth after 100-500 FODO cells



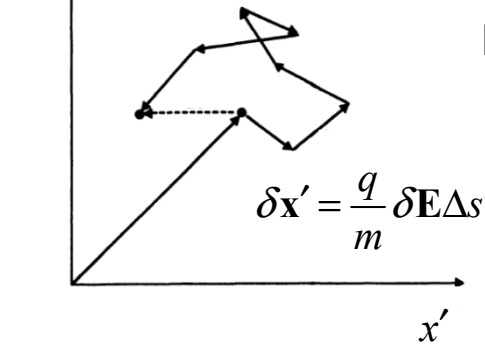
PATRIC (2D PIC):

FODO vs. CF



Stochastic or grid heating (,inelastic collisions‘)

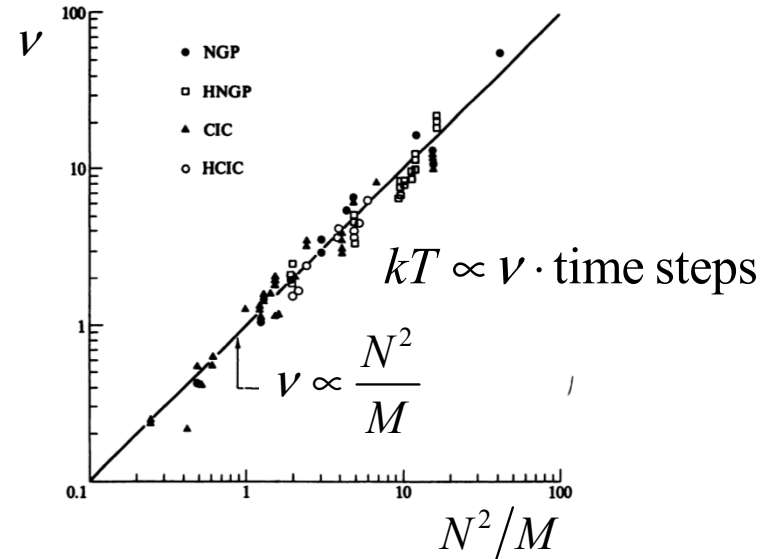
Random walk in grid
induced error fields:



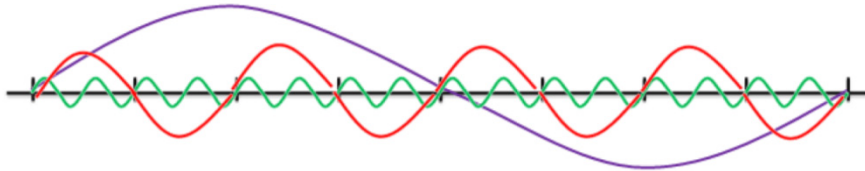
$$kT \approx n \frac{1}{2} \frac{q^2}{m} (\delta \mathbf{E})^2 (\Delta s)^2$$

Numerical heating due to grid !

$$D > \nu \frac{k_B T}{m}$$



$\delta \mathbf{E}$ Random grid induced electrical error field:
rounding, time step, finite differences,...



A non-conservative force (-> heating)
is introduced by **grid undersampling**.

-> Standard PIC integrators are non-symplectic

,Noise is not bad, the grid (undersampling) is‘

Stochastic heating in computer plasmas:

[1] Hockney, *Measurement of collision and heating times in a 2D computer plasma*, J. Comput. Phys. (1971)

Application to computer beams:

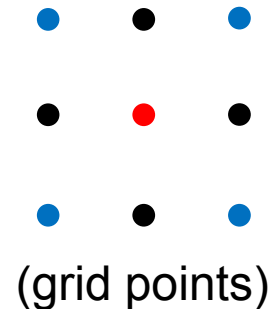
[2] Kesting, Franchetti, *Propagation of numerical noise in PIC*, PRAB (2015)

Digital filters for PIC

- + Reduce noise at short wavelengths (smoothing of the charge density).
- + Relatively fast.
- Can lead to non-physical results and numerical instabilities
- (Non-) Energy conservation (e.g. damping of modes)

$$\rho_{i,k} = \frac{M \rho_{i,k} + S(\text{side terms}) + K(\text{corner terms})}{M + 4(S + K)}$$

(M, S, K) = Middle, side and corner weights



Examples:

Hollow 8-point filter (0,1,1): Studied in [2]. Has negative spectral representation.
 Binomial 9-point filter (4,2,1): Proposed for PIC codes in [1].

Digital filters:

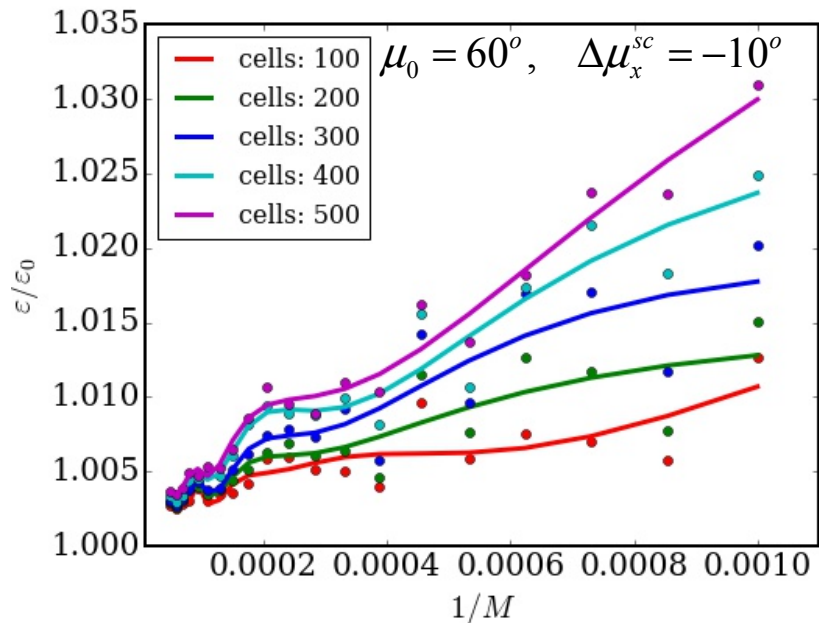
- [1] Birdsall, Langdon, *Plasma Physics via Computer Simulation* (2004)
- [2] Hockney, Eastwood, *Computer Simulation using Particles* (1988)
- [3] Verboncoeur, *PIC for plasmas: Review and Advances*, Plasma Phys. Control. Fusion (2005)
- [4] Vay et al., *Numerical methods for instability mitigation in ...*, J. Comput. Phys. (2011)

Wavelet denoising:

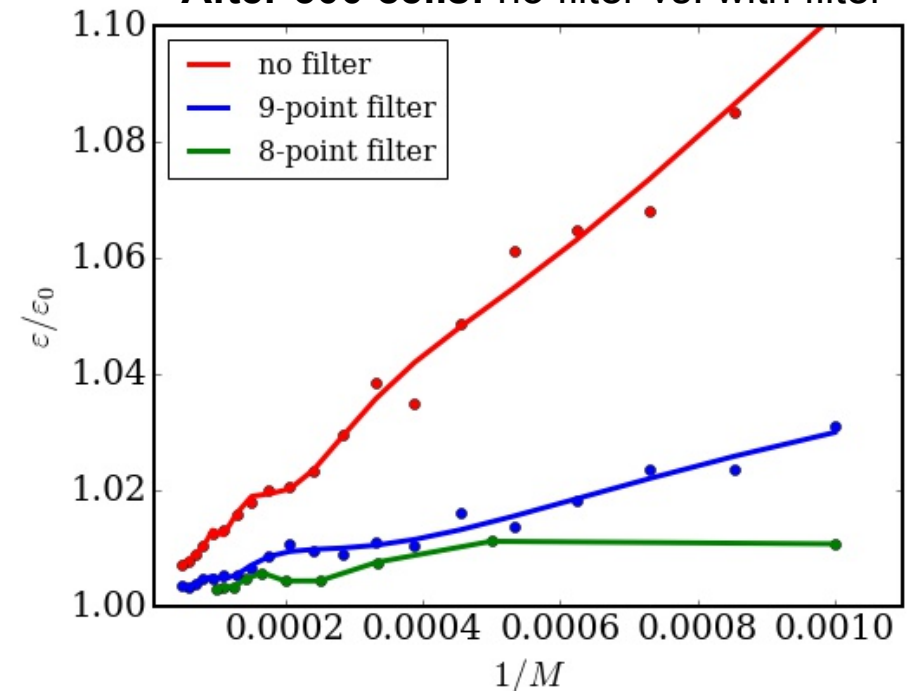
- [4] Gassama, Sonnendrücker, *Wavelet denoising for postprocessing of 2D PIC codes*, ESAIM (2007)
- [5] Terzic et al., *PIC simulations with a wavelet based Poisson solver*, PRAB (2007)

Example results for FODO channel with filters

PATRIC 2D: Numerical emittance growth in a FODO channel after 100-500 cells with 9-point filter



After 500 cells: no filter vs. with filter



With digital filters:

- > Emittance growth rate is slower (approx. factor 10).
- > However, emittance growth is not the only benchmark -> resonances,...

‘Symplectic Particle-In-Cell’

Conventional PIC: Total momentum is conserved, but not energy (grid heating).

‘Energy conserving’ PIC [3]: Momentum is not conserved (self-force on the grid).

-> PIC codes are usually not symplectic

An ideal beam particle integrator with space charge would conserve:

Phase space (Symplectic), Emittance, Entropy, Energy (only for constant focusing)

‘Multi-symplectic’ integrators offer bounded variations of those constants.

Starting point is usually the Lagrangian (here for a 2D beam with space charge):

$$L = \int dx_0 dy_0 dx'_0 dy'_0 f(x_0, y_0, x'_0, y'_0) \left(\frac{1}{2} ((x')^2 + (y')^2) - \frac{1}{2} \kappa (x^2 - y^2) - \frac{q}{E_0 \gamma_0^2 \beta_0^2} \phi(x) \right) - \frac{\dot{\phi}_0}{E_0 \gamma_0^2 \beta_0^2} \int dx dy \nabla \phi \cdot \nabla \phi$$

Equations of motion (for particles and fields) from Euler-Lagrange equations:

$$\frac{\partial L}{\partial \xi} - \frac{d}{ds} \frac{\partial L}{\partial \xi'} = 0 \quad \Rightarrow \quad x'' + \kappa(s)x = -\frac{q}{E_0 \gamma_0^2 \beta_0^2} \frac{\partial \phi}{\partial x} \quad y'' - \kappa(s)x = -\frac{q}{E_0 \gamma_0^2 \beta_0^2} \frac{\partial \phi}{\partial y} \quad \dot{\phi}_0 \Delta \phi = -\rho$$

[1] Webb, *A spectral canonical electrostatic algorithm*, Plasma Phys. Control. Fusion (2016)

[2] Shadwick et al., *Variational formulation of macro-particle plasma simulation algo.*, PoP (2014)

[3] Langdon, *„Energy-Conserving“ plasma simulation algorithms*, J. Comput. Phys. (1973)

[4] Qin, et al., *Canonical symplectic PIC method for long term simulations..*, arxiv 1503.08334 (2015)

Spectral (grid-less) computer beam Lagrangian

Beam macro-particle distribution:

$$f(\mathbf{x}, \mathbf{x}') = w \sum_{j=0}^M S(\mathbf{x} - \mathbf{x}_j) \delta(\mathbf{x}' - \mathbf{x}'_j)$$

Spectral space charge potential

$$\phi(\mathbf{x}) = \sum_{\mathbf{k}} \phi_{\mathbf{k}} \exp(i\mathbf{k} \cdot \mathbf{x})$$

Resulting discrete Lagrangian: $L_D = W_k - W_{AG} - W_{sc} - W_{\text{field}}$

$$W_k = w \sum_{j=0}^M \frac{1}{2} (\mathbf{x}'_j)^2$$

(kinetic)

$$W_{sc} = w \frac{q}{E_0 \gamma_0^2 \beta_0^2} \sum_{j=0}^M \sum_{\mathbf{k}} S(\mathbf{k}) \phi_{\mathbf{k}} \exp(i\mathbf{k} \cdot \mathbf{x}_j)$$

(space charge potential)

$$W_{\text{field}} = \frac{1}{E_0 \gamma_0^2 \beta_0^2} \sum_{\mathbf{k}} \sum_{\mathbf{k}'} k k' \phi_{\mathbf{k}} \phi_{\mathbf{k}'}$$

(space charge field energy)

Equations for macro-particles and spectra field) from Euler-Lagrange equations:

$$\frac{\partial L_D}{\partial \xi_j} - \frac{d}{ds} \frac{\partial L_D}{\partial \xi'_j} = 0 \quad \Rightarrow \quad \mathbf{x}''_j + \kappa(s) \mathbf{x}_j = - \frac{iq}{E_0 \gamma_0^2 \beta_0^2} \sum_{\mathbf{k}} \mathbf{k} \phi_{\mathbf{k}} S(\mathbf{k}) e^{i\mathbf{k} \cdot \mathbf{x}_j} \quad \epsilon_0 k^2 \phi_{\mathbf{k}} = \rho_{\mathbf{k}}$$

Spectral (grid-less) PIC for plasmas:

- [1] Webb, *A spectral canonical electrostatic algorithm*, Plasma Phys. Control. Fusion (2016)
- [2] Decyk, *Spectral PIC codes*, ISSS (2011)
- [3] Huang et al., *Grid instability and spectral fidelity of ES PIC codes*, arxiv:1508.03360 (2016)

Spectral algorithm for 2D beams

Spectral charge density: $\rho_k = wS(\mathbf{k}) \sum_j e^{-ik \cdot x_j}$

Space charge potential: $\epsilon_0 k^2 \phi_k = \rho_k$

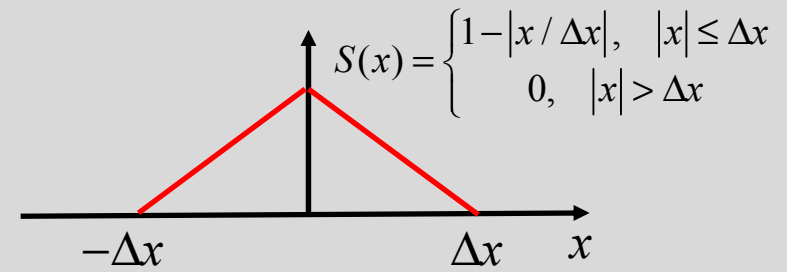
Space charge kick:

$$\Delta x'_n = -\Delta s \frac{iq}{E_0 \gamma_0^2 \beta_0^2} \sum_k k_x \phi_k S^*(\mathbf{k}) e^{ik_x \cdot x_j^n}$$

Time step (Δs):

$$\begin{pmatrix} x_j \\ x'_j \\ y_j \\ y'_j \end{pmatrix}_{n+1} = M(s_n, s_{n+1}) \begin{pmatrix} x_j \\ x'_j + \Delta x'_j \\ y_j \\ y'_j + \Delta y'_j \end{pmatrix}_n$$

'Tent' particles shape:



$$S(k) = \frac{1}{\sqrt{2\pi}} \int dx S(x) \exp(-ikx)$$

2D spectral particle shape:

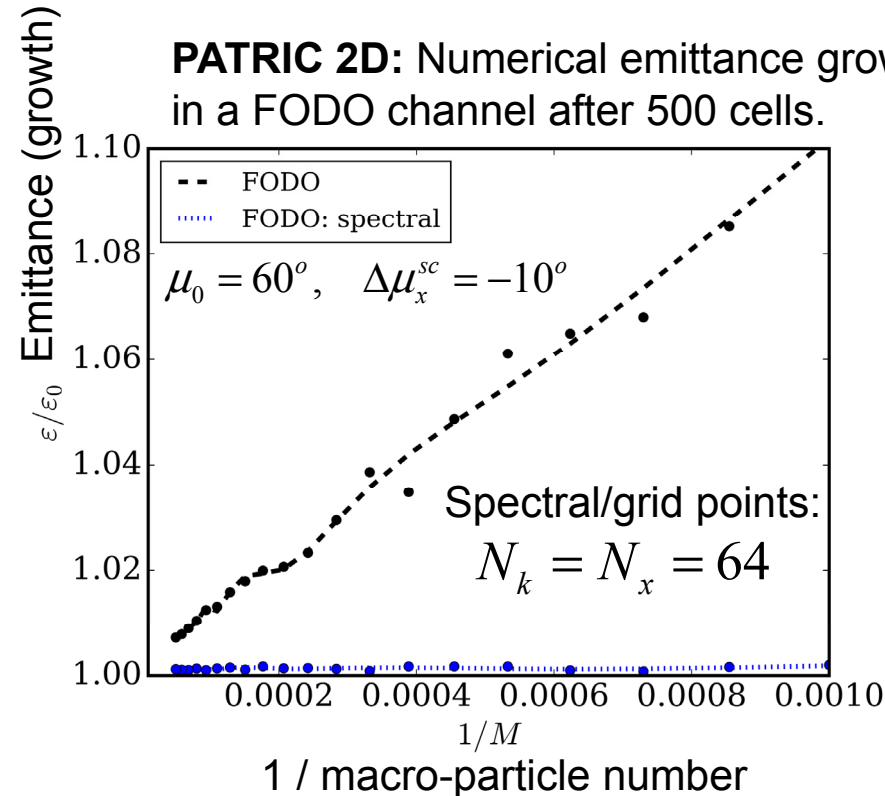
$$\Rightarrow S(\mathbf{k}) = \text{sinc}^2(\Delta x k_x / 2) \text{sinc}^2(\Delta y k_y / 2)$$

Algorithm derived from a grid-less macro-particle Lagrangian.

We expect the system to follow: $D \approx v \frac{k_B T}{m}$

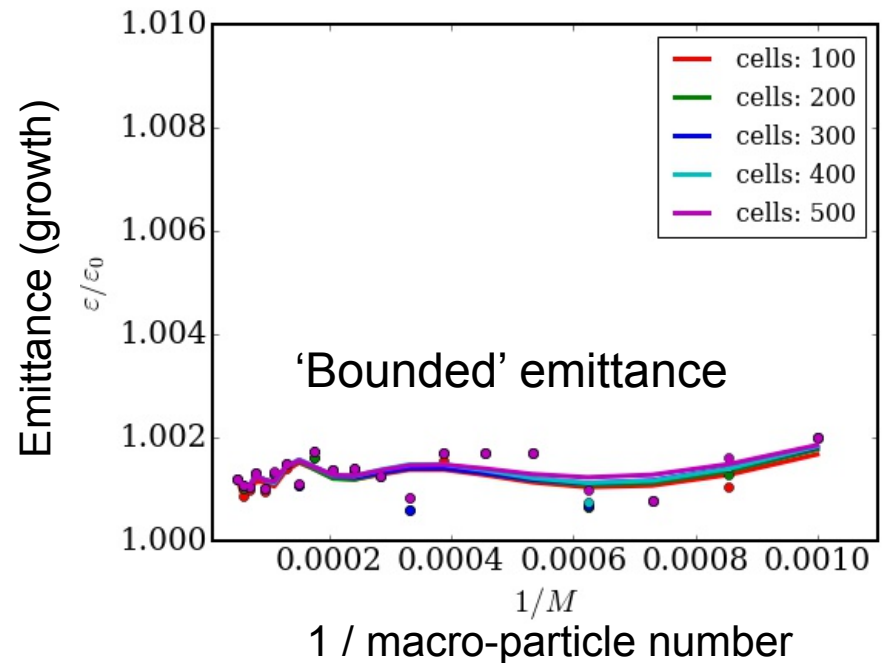
Tests: Spectral PIC vs. Conventional PIC

PATRIC 2D: Numerical emittance growth in a FODO channel after 500 cells.



Noise still present (in k-space), but no emittance growth.

Spectral PIC only.



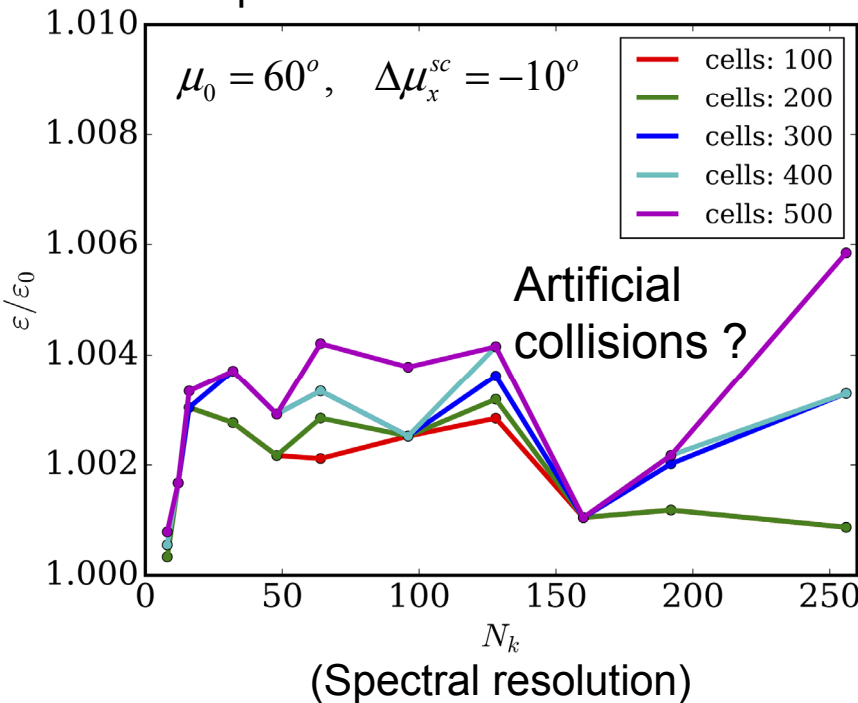
However, spectral PIC is (still) much slower:

$$T \propto \alpha \cdot M \cdot \text{cells} \quad \frac{\alpha_{\text{spectral}}}{\alpha_{\text{grid}}} \approx 20 - 50$$

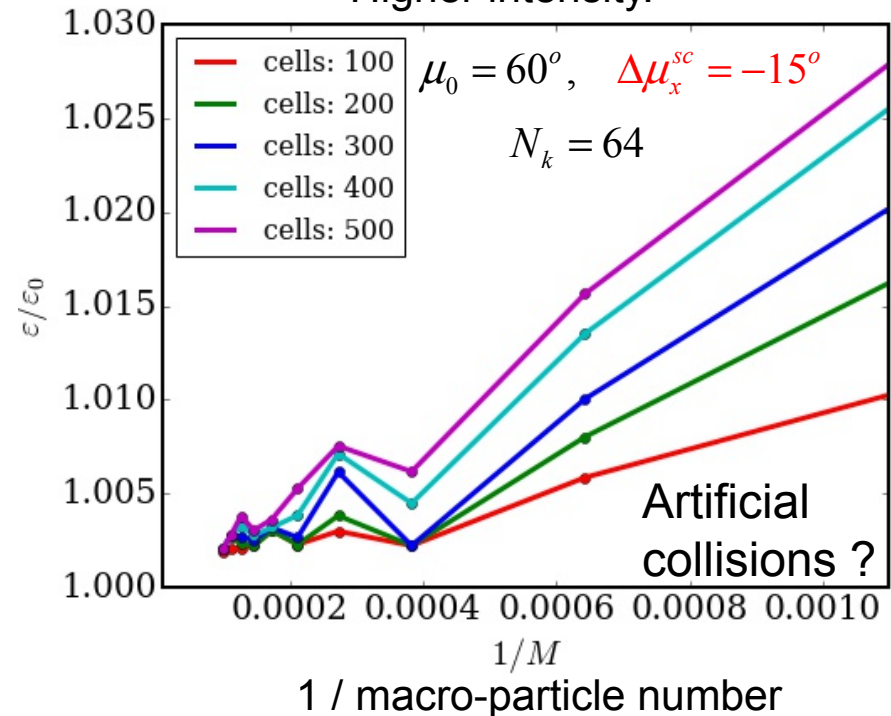
Remark: For a python/numpy/cython implementation and on a Mac Pro. Process running on one core.

Spectral PIC: More results for FODO channel

Numerical emittance growth in a FODO channel after 100-500 cells.
Spectral PIC with M=1000.



Higher intensity.



Because 'symplectic' PIC solves a macro-particle model, collisions should be present and the emittance should growth in AG focusing according to (still to be shown) :

$$D = \nu \frac{k_B T}{m} \Rightarrow \frac{1}{\epsilon} \frac{d\epsilon}{dt} = \frac{1}{2} k_B \nu \frac{(T_x - T_y)^2}{T_x T_y}$$

Conclusions and Outlook

Numerical emittance growth is a concern for long-term PIC particle tracking simulations with space charge in synchrotrons.

Two sources of numerical emittance growth:

- 1) Artificial collisions (or 'numerical IBS'): Emittance growth in AG focusing
- 2) Stochastic or grid heating ('inelastic collisions') -> dominant in conventional PIC

Collision rate: $\nu \propto \frac{N^2}{M}$ $D > \nu \frac{k_B T}{m}$ (diffusion, not balanced by friction -> heating)

Counter measures to decrease the growth rate: Larger M, particle shapes, filters,

'Symplectic' PIC:

- + Still collisional (we don't solve the Vlasov-Poisson equations) !
- + Noise/collisions do not cause emittance growth -> Bounded emittance $D \approx \nu \frac{k_B T}{m}$
- Usually slower than conventional PIC.

To do:

- > optimized implementations, e.g. using a properly chosen S(x) one a grid.
- > check other benchmarks: Resonances, Landau damping and echoes,.....
- > compare with other approaches, like symplectic Vlasov-Poisson solvers